### herramientas de

HPC para

física de colisiones y reacciones químicas

### herramientas de

HPC para

física de colisiones y

reacciones químicas

### $A + BC \longrightarrow A + B + C$

 $A + BC \longrightarrow A + B + C$ 

 $e + H \longrightarrow e + H^+ + e ionización$ 

 $A + BC \longrightarrow A + B + C$  $e + H \longrightarrow e + H^+ + e$ 

 $H^+ + H \longrightarrow H^+ + H^+ + e$  ionización

 $A + BC \longrightarrow A + B + C$  $e + H \longrightarrow e + H^+ + e$  $H^+ + H^- \longrightarrow H^+ + H^+ + e$  $Li + Li_2 \longrightarrow Li + Li + Li disociación$ 

inducida por colisión

 $A + BC \longrightarrow A + B + C$  $e + H \longrightarrow e + H^+ + e$  $|H^+ + H| \longrightarrow |H^+ + |H^+ + e$  $Li + Li_2 \longrightarrow Li + Li + Li$  $hv + He \longrightarrow e + He^{++} + e$ 

### A + B + C $e + H^{+} + e$ $H^{+} + H^{+} + e$ Li + Li + Li $e + He^{++} + e$ Doble Fotoionización



problema de 3 cuerpos cuántico

ecuación de schrödinger independiente del tiempo

### 3 cuerpos x 3 coord. c/u - 3 coord. CM = 6 coord

### esféricas

### coordenadas esféricas (2 x 3)



### coordenadas esféricas radiales (x2)



### funciones sturmianas generalizadas





dominio radial finito (2D) + C.C.

ecuación diferencial de segundo orden en derivadas parciales con condiciones de contorno complejas



dominio radial finito (2D) + C.C.

### discretización en diferencias finitas

### cálculo de la matriz H roblema de álgebra

problema de álgebra lineal numérica ralo



### método espectral con funciones sturmianas

# cálculo de la matriz H problema de álgebra

lineal numérica

### (e,3e) alta energía

### ionización simple ionización - excitación



### doble ionización



(e,3e) alta energía



m. ambrosio et al, enviado, 2014

UTE OF PHYSICS PUBLISHING J. Phys. A: Math. Gen. 36 (2003) 8443-8462 JOURNAL OF PHYSICS A: MATHEMATICAL AND GENERAL PII: S0305-4470(03)62313-9

### Two-body Coulomb wavefunctions as kernel for alternative integral transformations

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### Abstract

In this paper we investigate different representations of an arbitrary function in terms of two-body Coulomb eigenfunctions. We discuss the standard energy basis in spherical and parabolic coordinates with the purpose of remarking explicitly that two additional parameters appear both in the Schrödinge equation and in the wavefunctions: the charge and the angular momentum. We introduce the charge and generalized angular momentum Sturmian function representations, which result when the charge or the angular momentum is used as the eigenvalue in the Coulomb Schrödinger equation, respectively We present the connection between the generalized angular momentum representation and the Kontorovich–Lebedev transform. Finally, we extend the angular momentum representation to six dimensions, which is suitable for further applications in the three-body Coulomb problem.

PACS number: 03.65.-w

The standard way of theoretically studying the properties of a wide variety of atomic and molecular systems is by using functional basis set for representing the physical magnitudes involved in the phenomena under analysis. According to the mathematical structure of quantum mechanics, each physical magnitude has a Hermitian operator associated and then its eigenfunctions become the natural basis set to be used. The most common basis set uses the *energy eigenfunctions*, i.e., the eigenfunctions of the Hamiltonian. However, in some cases the energy eigenfunctions are not the most convenient ones and alternative basis sets are necessary. 8443

0305-4470/03/318443+20\$30.00 © 2003 IOP Publishing Ltd Printed in the UK

### $L^2$ Discretization of Sturmian Wave **Functions for Coulomb-like Potentials**

### A. L. FRAPICCINI,<sup>1,4</sup> V. Y. GONZALEZ,<sup>2,4</sup> J. M. RANDAZZO,<sup>1,4</sup> F. D. COLAVECCHIA,<sup>3,4</sup> G. GASANEO<sup>1,4</sup>

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Received 15 June 2006; accepted 7 August 2006 Published online 5 October 2006 in Wiley InterScience (www.interscience.wiley.com). DOI 10.1002/qua.21220

**ABSTRACT:** In this work we introduce a method to construct Sturmian functions for general interaction potentials in two-body problems. We expand these Sturmians on a finite  $L^2$  space, using N Laguerre basis functions to obtain a discrete set of eigenvalues for positive and negative energies. Orthogonality and closure relations are thus rewritten for these expansions; completeness is achieved through increasing the basis size. We apply the method to the Coulomb and Herman and Skillman potential. We study the behavior of the functions obtained and their convergence for an overall range of energies. The Sturmian functions are applied to solve the Schrödinger equation for an active electron in a He-like system. © 2006 Wiley Periodicals, Inc. Int J Quantum Chem 107: 832-844, 2007

Key words: Sturmian functions; Coulomb potential; Coulomb screened potential; Pollaczek polynomials

### 1. Introduction

- he theoretical study of a wide range of atomic and molecular properties can be carried out
- Correspondence to: A. L. Frapiccini; e-mail: afrapic@uns.edu.ar Contract grant sponsor: ANPCYT (Argentina). Contract grant number: PICTR 2003/00437. Contract grant sponsor: Universidad Nacional del Sur
- (Argentina). Contract grant number: PGI 24/F027.
- Contract grant sponsor: CONICET. Contract grant number: PIP 5595.

through the representation of a particular physical magnitude by a functional basis set. Such functions are usually generated by solving the Schrödinger equation for the problem under study, or an approx imated problem similar to the one of interest. The eigenfunctions obtained are those associated to the energy eigenvalues. However, in some cases, these functions are not suitable to a given problem and other basis sets are necessary or more convenient. e.g., for convergency reasons. As an alternative to energy eigenfunctions one could find solutions to the same wave equation treating the energy as a fixed

CHAPTER SEVEN

### **Three-Body Coulomb Problems** with Generalized Sturmian Functions

### G. Gasaneo<sup>a</sup>, L.U. Ancarani<sup>b</sup>, D.M. Mitnik<sup>c</sup>, J.M. Randazzo<sup>d</sup>,

A.L. Frapiccini<sup>a</sup> and F.D. Colavecchia<sup>d</sup> <sup>a</sup>Departamento de Física, Universidad Nacional del Sur and IFISUR-CONICET, Bahía Blanca, Argentina <sup>b</sup>Théorie, Modélisation, Simulation, SRSMC, UMR CNRS 7565, Université de Lorraine, 57078 Metz, <sup>b</sup>Théorie, Modélisation, Simulation, SRSMC, UMR CNRS 7565, Université de Lorraine, 57078 Metz, <sup>c</sup>Instituto de Astronomía y Física del Espacio (IAFE), CONICET and Departamento de Física, FCEyN,

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- 1. Introduction
- 2. Generalized Sturmian functions 2.1 Definitions
- 2.2 Bound states
- 2.3 Scattering states
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- 3.2 Two-electron atoms
- 3.3 Finite mass exotic and molecular systems
- 3.4 Confined atoms
- 4. Three-body problems: scattering states
- 4.1 Introduction
- 4.2 Asymptotic behaviors

4.3 Driven equation for three-body scattering problems
4.4 Solving the driven equation with GSF (spherical coordinates)
Insights from the zero-angular-momentum wave in single and double ionization
4.5 Solving the system content of the system of t 5. Three-body scattering states applications



0953-4075/10/101001+05\$30.00

[7–14]. Introduction of electronic correlation in the OUBLOW AND States can be performed in different ways. For example, one can explicitly include the dependence in the interelectronic coordinate, using relative interparticle We extended these works, introducing a systematic way to generate Sturmians for any physically sound potential with different types of boundary conditions [30–32].

Article history Received 7 July 2010 Received in revised form 17 Decem Accepted 24 January 2011 Available online 31 January 2011

Keywords: Generalized Sturmian functions Atomic spectra Continuum spectra

The computational techniques needed to generate a two-body Generalized Sturmian basis are described. These basis are obtained as a solution of the Schrödinger equation, with two-point boundary conditions. This equation includes two central potentials: A general auxiliary potential and a short-range generating potential. The auxiliary potential is, in general, long-range and it determines the asymptotic behavior of all the basis elements. The short-range generating potential rules the dynamics of the inner region. The energy is considered a fixed parameter, while the eigenvalues are the generalized charges. Although the finite differences scheme leads to a generalized eigenvalue matrix system, it cannot be solved by standard computational linear algebra packages. Therefore, we developed computational routines to calculate the basis with high accuracy and low computational time. The precise charge eigenvalues with more than 12 significant figures along with the corresponding wave functions can be computed on a single processor within seconds.

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1. Introductio

limiting first zero. The bound-hydrogenic functions are also widely

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A General Method to Obtain Sturmian

Functions for Continuum and Bound

State Problems With Coulomb

A. L<sup>2</sup>D1 FRAPICCINI,<sup>1,2</sup> J. M. RANDAZZO,<sup>1,2</sup> G. GASANEO,<sup>3</sup> F. D<u>2</u>QØLAVECCHIA<sup>1,2</sup>

Discrete Sets of Many-Body Sturmians

J. M. RANDAZZO,<sup>1,2</sup> A. L. FRAPICCINI,<sup>1,2</sup> F. D. COLAVECCHIA<sup>1</sup>

3. GASANEO<sup>-</sup> Centro Atómico Bariloche and CONICET, 8400 Bariloche, Rio Negro, Argentina In A. Frisia, Universidad Nacional del Sur and CONICET, 8000 Bahia Blanco

Received 12 December 2007; accepted 26 February 2008 Published online 14 August 2008 in Wiley InterScience (1

### Key words: Sturmians functions: generalized Stu 1. Introduction states of the t where the posi-the spherical or relative to the $\begin{array}{l} 154\\ \text{ the configuration interaction (CI) method max}\\ \text{ been Widely used to deal with two-electron}\\ \text{ atomic bound states. Roughly, it is based on the}\\ 157 \end{array}$ relative to the nuccess of international equation in sion, one arrives to a two-dimensional equation in the radial electronic coordinates, whose solution is approximated as a superposition of two-electron configurations. Other methods in which the basis expands the behavior of the wave function in all the interparticle distances are more rapidly convergen (see [1, 2] and references therein), but generally the calculation is very time consuming, and the algebr Correspondence to: J. M. Randarzoy e-mail: randarzo@cab.cnea. var 163 Contract grant sponsor: ANICYT (Argentina). Contract grant sponsor: CONICET. Contract grant sponsor: CONICET. Contract grant sponsor: Universidad Nacional del Sur

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### Contents lists available at ScienceDirect

### **Computer Physics Communications**

Computer Physics Communications 185 (2014) 1955–1964

journal homepage: www.elsevier.com/locate/cpc



### Accelerating spectral atomic and molecular collisions methods with graphics processing units

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functions. In the SFs for negative energies nake use of this method to	ARTICLE INFO	ABSTRACT					
Sturmian problem. We we CSF states, and show n energy-dependent This expansion in the sets used previously to infinity. © 2008 Wiley >velectron atoms	Article history: Received 10 January 2014 Received in revised form 14 March 2014 Accepted 26 March 2014 Available online 3 April 2014	We present a computation method to accelerate the calculation of the Hamiltonian of a three-body time independent Schrödinger equation for collisions. The Hamiltonian is constructed with one dimensional (basis overlaps) and two dimensional (interparticle interaction) integrals that are mapped into a computational grid in a Graphics Processing Unit (GPU). We illustrate the method for the case of an electron impact single ionization of a two electron atom. This proposal makes use of a Generalized Sturmian Basis set for each electron, which are obtained numerically on a quadrature grid that is used					
the wave function in a set of eigen- otal angular momentum operator $L^2$ , tions of the electrons are determined by coordinates which locate the electrons e nucleus. By means of such expan- res to a two-dimensional equation in	Keywords: Ionization Spectral methods Sturmian functions GPU computing	to compute the integrals in the GPU. The optimal computation is more than twenty times faster in the GPU than the calculation in CPU. The method can be easily scaled to computers with several Graphics Processing Units or clusters.					

### 1. Getting started

The properties of atoms and molecules can be determined by their wave function, which is obtained as the solution of the Schrödinger equation for non-relativistic energies. There are many methods to solve this second order, partial differential equation, from simple approximations that can give a hint of the physics, to the complete ab-initio numerical solution to predict physical quantities with a high degree of precision. When the wave function of an atomic or molecular system is written in a basis expansion, the Schrödinger equation transforms into a linear problem described by a dense matrix [1]. To obtain this matrix, it is necessary to compute integrals between two elements of the basis (called overlaps) and integrals of the interparticle potentials, that usually involve four basis functions.

The details of the calculation of each element of this matrix (called the Hamiltonian matrix) greatly depend on two factors. First, the system of coordinates chosen to represent the positions of the particles and their interactions, and second, the election of the basis set for each particle. These choices are determined not only by the number or type of particles involved, but also for the kind of experiment and/or physical property under investigation.

atomic systems, but not to analyze problems where particles can be spread out through the space, like ionization processes. Quantum chemistry calculations usually employ Gaussian functions to obtain properties of molecules, and take great advantage from the fact the systems under scrutiny are bound states [3–5]. Unlike these calculations of the quantum chemistry arena, colli-

sional problems deal with particles that can be far from each other In fact, the most important feature of a collision, the cross sections, are defined in those regions, assuming that the interactions among the particles are no longer effective and that they are far away from where the collision took place [6]. Therefore, the basis in collision problems should be able to accurately expand the full wave function for large interparticle distances. Moreover, if charged parti-cles are present, the basis should take into account the long-range asymptotic behavior of Coulomb fields [7,8]. Finally, possibly the main shortcoming is that it is not possible to use simple analytic basis sets for collision problems. All these factors pose several challenges to perform a numerically accurate calculation of wave func-tions and cross sections in atomic and molecular collisions [9].

In the last few years, the computer hardware and software have been moving fast to an heterogeneous world [10]. In the desktop market, this meant going from a simple one-core desktop er to an aggregate of one to several multicore CPUs with prresponding accelerators. This also replicates in High uting clusters, such as the Titan supercomp nal Laboratory [11]. Nowadays it is not possible ntific code thinking only about the pure speed (i.e., floating points operations per second). In-pre communication layers, memory hierarchies,

<ul> <li>and Province Barnovne, 9400 S. C. te Barnovne, Ko Vegro, Avgentinad es Ortificos y Técnicas, CONCET, Argentina ad Nacional del Sur, 8000 Bahía Blanca, Buenos Aires, Argentina</li> <li>2 Interna 2008 ey currS vce (usor or ner mcs 'ley.com).</li> </ul>	(exponentials times powers of the elect to the unit of the a basis of 2]. It is a basis of 2].	ronic coordinates relative of study both	tais their corresponding accelerators. This also Performance Computing clusters, such as the Ti at Oak Ridge National Laboratory [11]. Nowada to program a scientific code thinking only ab of the calculation (i.e., floating points operatio	rep tar ys out
	E-muil dudress: Havioce D.chea.gov.ar.	-	core and out-of-core communication layers, in	.en
investigate discretization schemes to represent and negative energies in the presence of a long range methods to obtain Sturmian functions for positive spansion of the radial wave function in a L <sup>2</sup> finite basis uces the discretization of the radial coordinate domain the Green function. We apply them to find the bound N-delectron atoms close to the critical charge. Both	http://dx.doi.org/10.1016/j.cpc.2014.03.026 0010-4655/© 2014 Elsevier B.V. All rights reserved	d.	Available online at www.sciencedirect.com	
d states near threshold, as well as continuum states			ScienceDirect	
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Few-Body Syst scattering; phase shift; critical <b>Dodg</b> t0.1007/s00601-014-0831-5		ELSEVIER	Journal of Electron Spectroscopy and Related Phenomena 161 (2007) 199-203	wv
		Stu	urmian functions in a $L^2$ basis: Critical nucle	ear

Contract grant sponsor: ANPCYT (Argentina). Contract grant numbers: PICTR 2003/0437 and PICT 04/20548. 4/2038. Contract grant sponsor: Universidad Nacional del Sur Irgentina).

Contract gran Tumber: PGL 24/038

International Journal of Quantum Chemistry, Vol 109, 125-134 (2009) © 2008 Wiley Perjodicals, Inc.

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Interactions

153ACT: In this article, we functions for both posit tunctions for both positi Coulomb potential. We explore tw energy. The first one involves the e set, whereas the second one introd of the Hamiltonian or alternatively

ates and scattering phase shift in methods are able to describe bou

Key words: Sturmian func

International Journal of Quantum Chemistry, Vol 110, 963–974 (2010) © 2009 Wiley Periodicals. Inc.

PHYSICAL REVIEW A 84, 052715 (2011)

A Generalized Sturmian Treatment of (e, 3e) Processes

Described as a Three-Body Coulomb Problem

M. J. Ambrosio · L. U. Ancarani · D. M. Mitnik

E.D. Colavecchia · G. Gasaner

The three-body breakup problem is of fundamental interest in atomic collisions theory. The simplest example, ionization of atomic hydrogen by electron impact, is theoretically de-scribed by the solution of the Schrödinger equation associated with two electrons moving in the field of a heavy nuclei con-sidered at rest. Analytical solutions are not known. However, cross sections for various energies have been measured in the laboratory for this process.

laboratory for this process. The main theoretical difficulties of the three-body fragmentation problems with Coulomb interactions are related to the very complicated form of the six-dimensional (correto the very complicated form of the six-admensional (corre-lated) wave function together with its asymptotic properties [1–5]. Many separable and nonseparable models have been introduced to obtain approximate values for the transition amplitudes [6,7], which are solutions of the Schrödinger equation in some asymptotic region where at least one of the particles is far away from the other two. However, the

Preprint method (FM) calculation performs the text of the second one is a finite element method (FM) calculation performs the second one is a finite element method (FM) calculation performs the second one is a finite element method (FM) calculation performs the second one is a finite element method (FM) calculation performs the second one is a finite element text of FAM text finite for a second one is a finite element and the field second text field finite domain (CM) and the field second text field field

Sur, Bahía Blanca, Argentina. E-mail address: afrapic@uns.edu.ar (A.L. Frapiccini).

nikov and Macek [8] obtained a discrete set of eigenvalues fo purely outgoing wave Sturmians, however this functions became unbounded as *r* increased. Rawitscher [9] was able to define a set of Sturmians with outgoing wave condition even in the case where a long range potential was present, showing that they constitute a discrete basis set with discrete eigenvalues.

domain. This basis set is therefore suitable for constructing the wave function of a given scattering problem for both long range Coulomb potentials or short range potentials. In Section 1 of this paper we present a brief review of the Sturmian theory and an analysis of the different asymptotic behavior according to the energy domain. In Section 2 we outline the general method to expand the two-particle Sturmian functions in terms of Laguerre-*type* basis, and obtain orthogonality and closure relations restricted to a finite subspace. Numerical results for a Coulomb well potential are shown for both negative and positive energies. In Section 3 we use the

with box bounders, head to be a large square domain. Al-though outgoing, flux insbulic texts or scoordinates in telated 197078 Metz, France the double continuing charinelist, the ubitained wave function is valid in the whole spatial domain where overlapping with other channels occus. Gainese'n M: the minore dominan (" behavior, when Departamento de l'isica, Universidad Nacional del sur, 8000 Bahía Blanca, Buenos Aires, Argentina fragmentation take place. Accurate scattering wave functions and breakup cross sections are obtained with the ECS method, although at the price of the p

(13). Besides, F. D. Colavecchia cutoff in some fit with the framework of this method, an artuncial cutoff in some fit with collisiones and marked the work. CONICET, Bariloche, Río Negro, Argentina Other theories such as the J-matrix approach employ Other theories such as the *J*-matrix approach empiri-spectral techniques to deal with the three-body proble

Abstract

Two particle Sturmian functions [M. Rotenberg, Ann. Phys., NY 19 (1962) 262; S.V. Khristenko, Theor. Math. Fiz. 22 (1975) 31 Theor. Math. Phys. 22, 21)] for a short range potentials are obtained by expanding the solution of the Schrödinger equation in a finite type basis. These functions are chosen to satisfy certain boundary conditions, such as regularity at the origin and the correct asymp ns [M. Rotenberg, Ann. Phys., NY 19 (1962) 262; S.V. Khristenko, Theor. Math. Fiz. 22 (1975) 31 (Engl. Transl. short range potentials are obtained by expanding the solution of the Schrödinger equation in a finite L<sup>2</sup>Laguerre-J. M. Randazzo,<sup>1,\*</sup> F. Buezas,<sup>2</sup> A. L. Frapiccini,<sup>1</sup> F. D. Colayecchia,<sup>1</sup> and Conferentials are obtained by expanding the solution of the schrödinger equation in a finite *L*<sup>2</sup> *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che motion france schrödinger equation in a finite L*<sup>2</sup> *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che motion france schrödinger equation in a finite L*<sup>2</sup> *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che motion france schrödinger equation in a finite L*<sup>2</sup> *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che motion france schrödinger equation in a finite L*<sup>2</sup> *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che motion france schrödinger equation in a finite L*<sup>2</sup> *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che schrödinger equation*, solution at the schrödinger equation in a finite *C*-1 *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che schrödinger equation*, and the correct asymptotic schrödinger equation in a finite *L*<sup>2</sup>. Straspic *C* - 1 *Departamento de Física*, *Universidad Nacional del Sur and CONICET*, 8000 *Barilo che schrödinger equation*, and *CONICET*, 8000 *Barilo che schrödinger equation*, and *Barilo che* 

charge for *N*-electron atoms

A.L. Frapiccini<sup>a,d,\*</sup>, G. Gasaneo<sup>a,d</sup>, F.D. Colavecchia<sup>b,d</sup>, D. Mitnik<sup>c,d</sup>

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Available online 21 February 2007

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bloque de L fijo





### bloque de L fijo











# Cada elemento de matriz H tiene

### n integrales 2D (repulsión)

2 integrales ID (overlaps) prod. escalar

prefix sum prod. escalar



2 He Helium 4.003

### L=0

### $\times$ |3 x |3 x 2

### 253 605 625 4

# $x | 3 pares (|_a,|_b) x n=2 x N_a = N_b = 35$ × 1225 × 1225

2DGB

colisión (e,3e)



### L=0-4 x 5 pares $(I_a, I_b)$ x n=2 x $N_a = N_b = 84$

### 5

## 12 446 784 000 2D 98

# x 5 x 5 x 2 x 7056 x 7056

GB



### método espectral

# cálculo de la matriz H

problema de álgebra lineal numérica denso

cálculo de la matriz H



# C + Fortran + CUDA 14000000 2D/s $25 \times$



### Accelerating spectral atomic and molecular collisions methods with graphics processing units

F.D. Colavecchia

### ARTICLE INFO

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COMPUTER PHYSICS COMMUNICATIONS

División Física Atómica, Molecular y Óptica, Centro Atómico Bariloche, and CONICET, 8400 S. C. de Bariloche, Río Negro, Argentina

### ABSTRACT

We present a computation method to accelerate the calculation of the Hamiltonian of a three-body time independent Schrödinger equation for collisions. The Hamiltonian is constructed with one dimensional (basis overlaps) and two dimensional (interparticle interaction) integrals that are mapped into a computational grid in a Graphics Processing Unit (GPU). We illustrate the method for the case of an electron impact single ionization of a two electron atom. This proposal makes use of a Generalized Sturmian Basis set for each electron, which are obtained numerically on a quadrature grid that is used to compute the integrals in the GPU. The optimal computation is more than twenty times faster in the GPU than the calculation in CPU. The method can be easily scaled to computers with several Graphics Processing Units or clusters.

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### problema de álgebra lineal numérica denso

### Fortran + MPI + ScaLapack

### d. mitnik

north carolina (EE UU) metz (Francia)

# cálculo de la matriz H

problema de álgebra lineal numérica denso

> producción ler semestre 2015

# C + Fortran + CUDA MPI Fortran + MPI + ScaLapack





### 4 cuerpos



-







### problema de álgebra lineal numérica denso

para

### una matriz que no se conoce a priori

pero cuyos elementos se saben calcular

### matrices latentes

### pict 2014 presentado

n wolovick c bederián e pilotta l biedma j m randazzo

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g gasaneo herramientas de

financiación

física de colisiones y reaccip Balseiro NVIDIA